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**Abstract.** Terahertz spectrometers and imaging systems are currently being evaluated as biomedical tools for skin burn assessment. These systems show promise, but due to their size and weight, they have restricted portability, and are impractical for military and battlefield settings where space is limited. In this study, we developed and tested the performance of a compact, light, and portable THz time-domain spectroscopy (THz-TDS) device. Optical properties were collected with this system from 0.1 to 1.6 THz for water, ethanol, and several *ex vivo* porcine tissues (muscle, adipose, skin). For all samples tested, we found that the index of refraction ( $n$ ) decreases with frequency, while the absorption coefficient ( $\mu_a$ ) increases with frequency. Muscle, adipose, and frozen/thawed skin samples exhibited comparable  $n$  values ranging between 2.5 and 2.0, whereas the  $n$  values for freshly harvested skin were roughly 40% lower. Additionally, we found that the freshly harvested samples exhibited higher  $\mu_a$  values than the frozen/thawed skin samples. Overall, for all liquids and tissues tested, we found that our system measured optical property values that were consistent with those reported in the literature. These results suggest that our compact THz spectrometer performed comparable to its larger counterparts, and therefore may be a useful and practical tool for skin health assessment. © 2011 Society of Photo-Optical Instrumentation Engineers (SPIE). [DOI: 10.1117/1.3570648]

Keywords: terahertz; THz; skin; terahertz time-domain spectroscopy; optical properties; compact.

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## 1 Introduction

The terahertz (THz) region of the electromagnetic (EM) spectrum is defined as frequencies ranging from 0.1 to 10 THz (1 THz =  $10^{12}$  Hz = 1 ps). Historically, few sources have been available to efficiently generate THz radiation. However, over the past few years, several technological advances have resulted in the unprecedented development of many new types of THz sources and components.<sup>1–5</sup> These technologies are now being used as tools for a plethora of basic science investigations.<sup>6–8</sup> In addition, they are increasingly being integrated into innovative sensing and imaging operational schemes, which are finding widespread use in a host of medical, military, and defense applications. For example, THz devices are now being tested at hospitals for early cancer and burn diagnosis,<sup>9–22</sup> at major airports for security screening purposes,<sup>22–24</sup> and at border patrol checkpoints for identification of explosives,<sup>3,25</sup> drugs,<sup>26,27</sup> liquids,<sup>7,8,28–30</sup> and biological macromolecules.<sup>31–34</sup>

The strength of many THz dermatological applications is derived from the fact that water, the primary constituent of skin tissues (~70% by volume), exhibits strong absorption at THz frequencies ( $\mu_a = 300 \text{ cm}^{-1}$  at 1.5 THz).<sup>28</sup> Due to the

strong optical absorption of water, THz waves can only penetrate a few hundred microns in biological tissues. Therefore, THz techniques are best suited for imaging the surface of biological tissues. However, the strong absorption of water is also beneficial because it provides for excellent soft-tissue contrast. Consequently, THz sensing and imaging approaches are very sensitive to small differences in the water concentration levels of tissues. Several recent studies have successfully exploited this contrast mechanism to differentiate healthy, dehydrated, burned, and diseased tissues.<sup>12–14</sup>

Water exhibits many unique chemical and physical properties that contribute to its strong interaction with THz radiation. One property is that water molecules are able to form tetrahedral arrangements of hydrogen bonds to neighboring molecules. This property gives water molecules the ability to actively engage in both inter- and intramolecular bonding. Due to these interactions, water molecules create an extensive dynamic hydrogen bond network, which behaves in a collective manner. Interestingly, the intermolecular stretching vibrations of this network, which are the origin of macroscopic water dynamics, occur on the picosecond timescale (5.6 THz).<sup>35–37</sup> In addition, due to the slow relaxation time of water molecules, large intermolecular bending vibrations also occur on the

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picosecond timescale (1.5 THz).<sup>35–37</sup> Water also functions as a universal solvent and is required for the proper folding and function of biological macromolecules (i.e., carbohydrates, proteins, and nucleic acids).<sup>38–42</sup> Many of these macromolecules also exhibit collective vibrational modes on the picosecond timescale.<sup>31,34,43–45</sup> Overall, water and many biological macromolecules exhibit dynamics that occur on the picosecond timescale.

Numerous basic science investigations have used THz spectroscopic techniques to characterize the frequency-dependent optical properties of water and biological molecules.<sup>7,8,28–30,32,46–49</sup> There are two advantages to using THz spectroscopic techniques rather than conventional microwave or Fourier transform far-infrared spectroscopic (FTFIR) approaches. The first advantage is that microwave spectroscopic approaches operate on the 100 ps timescale, whereas THz techniques operate on the subpicosecond to picosecond timescales. Since the vibrational motion of water and biomolecules occur on the picosecond timescale, THz approaches permit time-resolved investigation of these dynamics because they operate on the same timescale (1 ps = 1 THz). The second advantage of THz spectroscopic approaches is that they permit the measurement of amplitude and phase information, which are directly related to the absorption coefficient and index of refraction of the tissue. Thus, the complex permittivity of samples can be obtained without having to perform Kramers–Kronig analysis.<sup>50</sup>

THz spectroscopy is typically performed in either a reflection or transmission geometry. Transmission-based approaches have been used to examine the dynamics of water, biomolecules, and biological tissues. However, due to the strong absorption of water, a transmission-based measurement is typically performed on either lyophilized tissues or on very thin tissue slices ( $\leq 100\mu\text{m}$ ).<sup>15–17</sup> Lyophilization and tissue sectioning approaches make it feasible to conduct transmission-based measurements; however, these methods greatly affect the hydration level and morphology of the tissue. As a consequence, these changes can alter the optical properties that are collected, thereby reducing the accuracy of the measurements. Due to these factors, reflection geometries are preferred for both *in vivo* and *ex vivo* THz spectroscopic measurements of biological tissues.

Over the past decade, several reflection-based THz systems have been developed for the noninvasive evaluation of biological tissues.<sup>9–11,51–53</sup> These systems have been used to evaluate tissue hydration,<sup>15–17</sup> to diagnose skin cancers,<sup>9–14,53–56</sup> and to assess the depth and severity of skin burns.<sup>18–22</sup> These studies have demonstrated that the tissue's water concentration levels and its optical properties are the primary contrast mechanisms for differentiating healthy, burned, and diseased tissues.<sup>12–14</sup> These studies have also demonstrated that the optical properties of biological tissues depend on the organism (e.g., human, rat, porcine),<sup>10,11,52,53,57</sup> tissue type (e.g., skin, muscle, adipose, liver, heart),<sup>53,58,59</sup> method of tissue preparation (e.g., fresh versus frozen/thawed), and duration of tissue fixation.<sup>15,57</sup>

In recent years, many THz-TDS systems have become commercially available. These systems are useful for many applications, but due to their size (3000 to 50000 in<sup>3</sup> or 50 to 250 liters) and weight (55 to 220 pounds or 25 to 100 kg), they have restricted portability and are impractical for military and battlefield settings where space is limited. For example, in battle-

field hospitals or medical evacuation aircraft, and other Medevac transportation vehicles. In this study, we developed and tested the performance of a compact (180 in<sup>3</sup> or 2.95 liters), light (4.2 lbs or 1.9 kg), and portable THz-TDS device. This device is envisioned to be employed as a tool for the rapid assessment of skin wounds and/or burns (i.e., thermal, chemical, physical) in battlefield settings. To examine the performance of this device, we measured the optical properties for several liquids and biological tissues, and compared our results to those reported in the literature.

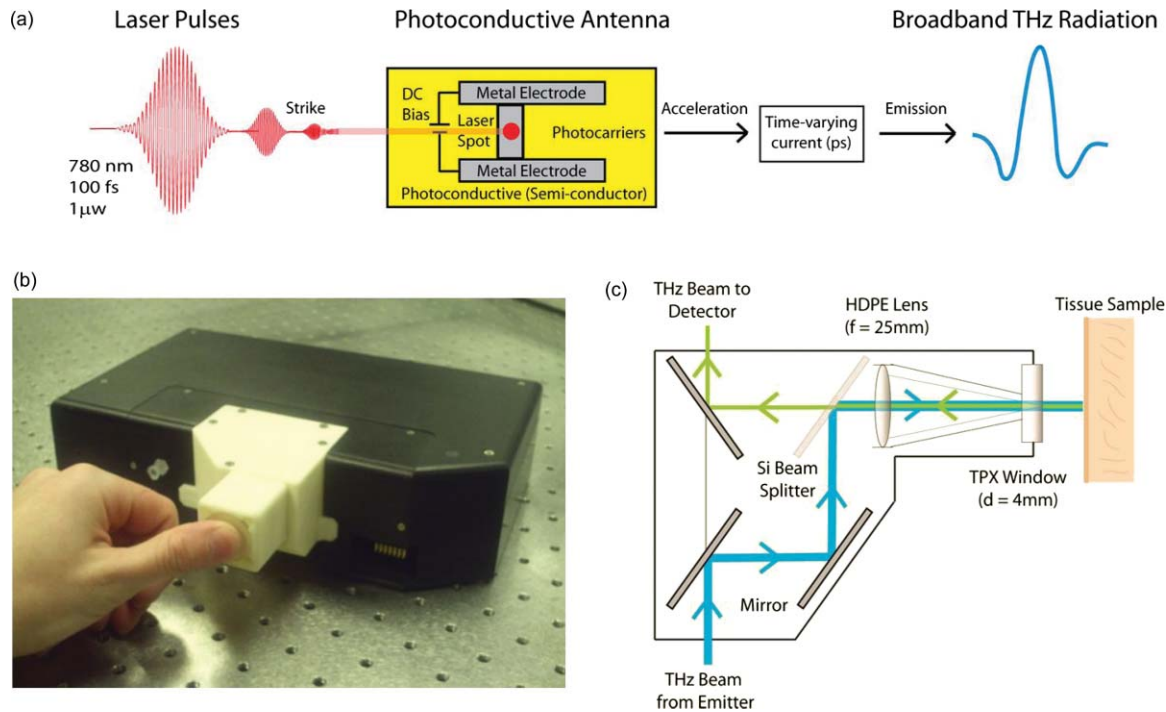
## 2 Materials and Methods

### 2.1 Liquid and Biological Sample Preparation

We measured the spectra for water, ethanol, and several porcine tissues using our compact THz-TDS device. All measurements were made at room temperature. Five or more measurements were made for each liquid, and each test was performed by placing a drop of liquid ( $\sim 50\mu\text{L}$ ) on the measurement window. Excised porcine tissues (skin, adipose tissue, and muscle) were collected and prepared in adherence to a Brooks Air Force Research Laboratory exempt Institutional Animal Care and Use Committee protocol. The thickness of all samples was greater than or equal to 1 cm. Spectra was measured on the following freshly excised porcine tissues: muscle, adipose, and skin. To examine the effect that tissue preparation has on the optical properties, we also collected and compared the spectra for freshly excised and frozen/thawed skin samples. The spectra of the fresh samples was measured on tissues less than one hour after excision. The frozen/thawed skin samples were prepared as follows: excised two days prior to experimentation, wrapped in moistened gauze, frozen at  $-20^\circ\text{C}$ , and thawed at room temperature 30 min prior to measurement. Seven to twenty-six measurements were made on each tissue sample, and samples were processed for conventional histological analysis using hematoxylin and eosin (H&E) stains.

### 2.2 Compact Terahertz Time Domain Spectrometer

In this study, we developed a modified THz-TDS system that permits reflection measurements for biological tissues. The system consists of a mini-Z<sup>TM</sup> THz-TDS system (Zomega Terahertz Corporation) and a custom-designed skin reflection unit [Figs. 1(a)–1(c)]. The system is compact (180 in<sup>3</sup> or 2.95 liters), light (4.2 lbs or 1.9 kg), inexpensive ( $\sim \$150\text{K}$ ), portable, and robust. It is interesting to note that this system is 15 to 55 times lighter and 15 to 80 times smaller than conventional commercially available systems. The unit has the following performance specifications: frequency range (0.1 to 2.5 THz), frequency resolution (5 GHz), dynamic range (70 dB at 0.1 THz to 50 dB at 1.6 THz). The mini-Z<sup>TM</sup> THz-TDS system generates broadband THz radiation using a photoconductive (PC) switch technique, originally pioneered by Auston and colleagues.<sup>60–63</sup> In brief, the PC switch technique functions by using a pulse from a femtosecond (fs) near-infrared laser to excite a biased PC antenna consisting of metallic striplines deposited on semiconductor materials. Upon interaction with the semiconductor material, the optical pulse generates photocarriers, which when accelerated using a DC bias, create a photocurrent within the PC antenna. These time-varying photocurrents, which occur in the subpicosecond range, emit broadband EM radiation at THz frequencies.



**Fig. 1** Compact terahertz time-domain spectrometer for the measurement of the optical properties of biological tissues. (a) Schematic illustration of the photoconductive switch used to generate THz radiation in our source. (b) Image of custom THz-TDS system consisting of a skin reflection unit and the Mini-Z™ spectrometer (Zomega Terahertz Corporation, LTD). (c) Schematic representation of principle optical elements used in the device.

Figure 1(a) contains a schematic illustration of the primary elements used in our system. The pump laser used in this unit is a diode laser with the following specifications:  $\lambda = 780$  nm,  $\tau_p = 100$  fs, and average power = 20 mW. Using a zinc telluride electro-optic crystal (1 mm thick), we obtained the time-varying electric fields of the THz reflected light which, upon performing a numerical Fourier transform, resulted in reflection spectra from 0.1 to 4.0 THz. The skin reflection unit contains three gold mirrors ( $R \approx 0.99$  at THz frequencies), a high resistivity (10K ohm-cm) Silicon beam splitter, a high-density polyethylene (HDPE) lens, and a polymethylpentene (TPX) window [Figs. 1(b) and 1(c)]. The HDPE lens has a diameter of 2.54 cm, a 25 mm focal length, and  $\mu_a = 0.1$  to  $5.0$   $\text{cm}^{-1}$  from 0.1 to 1.6 THz. The 4-mm thick TPX window has a diameter of 2.54 cm and  $\mu_a = 0.1$  to  $1.5$   $\text{cm}^{-1}$  from 0.1 to 1.6 THz. The window served two purposes: 1. to ensure samples were normal to the THz beam and 2. to accurately measure the phase of the reflected signal by monitoring the relative timing of the signal reflected from the sample compared to the reflection from the first surface of the TPX window.

### 2.3 THz TDS Source Performance

Figures 2(a) and 2(b) contain sample plots of the broadband THz pulses created with our THz-TDS device. A plot of the wave form of the broadband THz pulses (FWHM = 660 fs) in time domain is provided in Fig. 2(a), and a plot of the frequency power spectrum of the Fourier transform of the pulse is shown in Fig. 2(b).

### 2.4 Data Acquisition and Signal Processing

The raw data collected with our system is in the time domain and the optical properties ( $n$  and  $\mu_a$ ) of the samples are frequency

dependent. Therefore, we conducted a fast Fourier transform (FFT) to compute these values. To do this, we obtained the power reflectance ( $R$ ) and reflection phase ( $\theta$ ), and then took the complex Fourier transform of the sample measurements. The phase was accurately measured by comparing the reflection from the first and second interface of the TPX window. These signals were then compared to a reference consisting of a mirror pressed against the window. Each measurement was computed using an average of 250 waveforms acquired at a rate of 5 Hz. Given  $R$  and  $\theta$ , we then used Fresnel equations to calculate the  $n$  and the extinction coefficient ( $\kappa$ ) of the sample [see Eqs. (1) and (2)]

$$n = n_0 \frac{1 - R}{1 + R - 2\sqrt{R} \cos(\theta)}, \quad (1)$$

$$\kappa = n_0 \frac{-2\sqrt{R} \sin(\theta)}{1 + R - 2\sqrt{R} \cos(\theta)}, \quad (2)$$

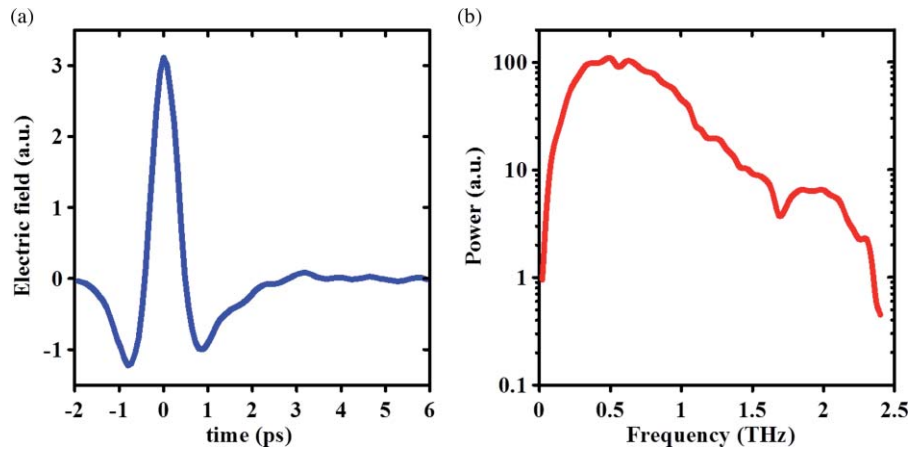
where  $n_0$  is the index of the TPX window. The absorption coefficient  $\mu_a$  was then calculated using the following relationship:  $\mu_a = 4\pi\kappa/\lambda$ , where  $\lambda$  is the free space wavelength. Then using Beer-Lambert Law, with the assumption that THz-liquid interactions are absorption-dominated, we then calculated the optical penetration depth ( $\delta = 1/\mu_a$ ).

## 3 Results and Discussion

### 3.1 System Calibration and Validation using Several Liquids

#### 3.1.1 Ethanol

The goal for the first set of experiments was to validate that our system provided comparable values to those reported in



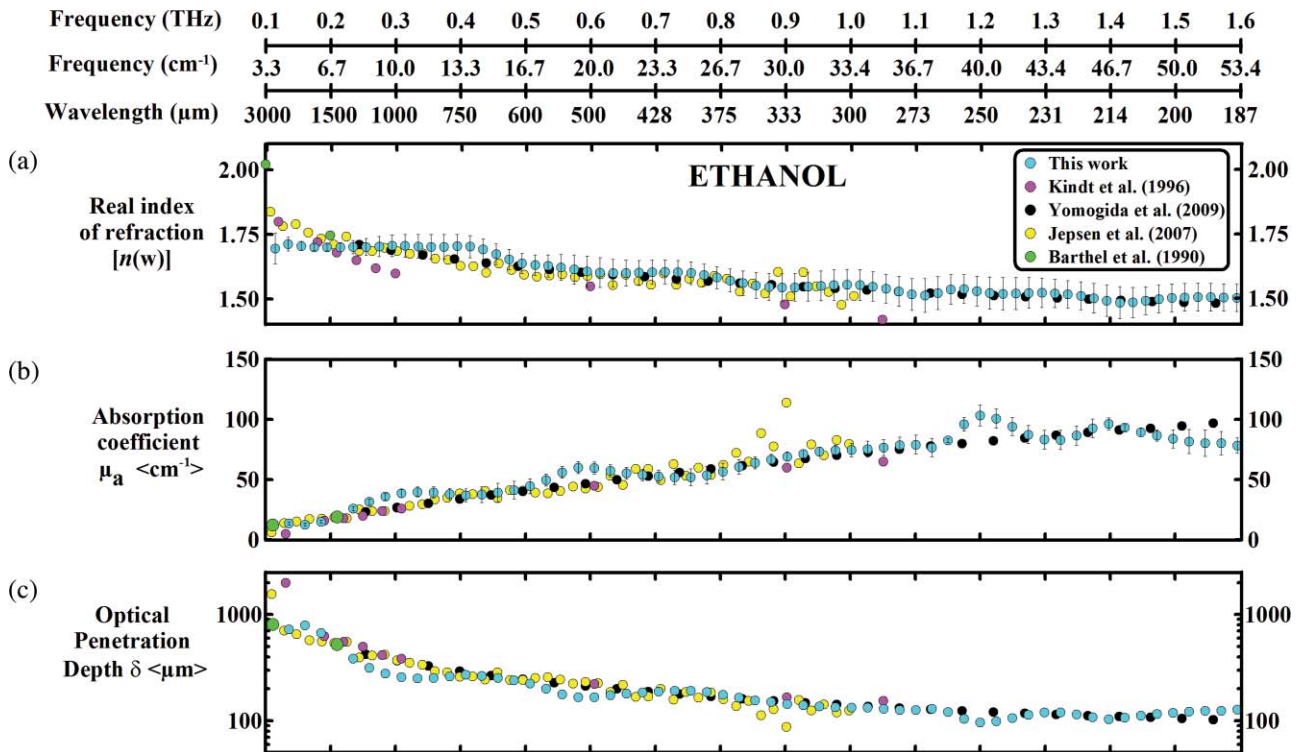
**Fig. 2** Performance of THz-TDS device. (a) Wave form of THz pulse (FWHM = 660 fs) measured in the time domain. (b) Frequency power spectrum of the Fourier transform of the pulse shown in (a).

the literature. Since the optical properties of ethanol are well-characterized at THz frequencies, we first used our system to measure the properties of ethanol. Figures 3(a) and 3(b) contain plots of the frequency-dependent optical properties for ethanol. We found that the index of refraction decreases with frequency, whereas the absorption coefficient increases with frequency. We also found that our collected values are in excellent agreement with previously published reports.<sup>28,32,64,65</sup> Figure 3(c) contains a plot of the optical penetration depth for ethanol. We found that at lower THz frequencies, ethanol has an optical penetration

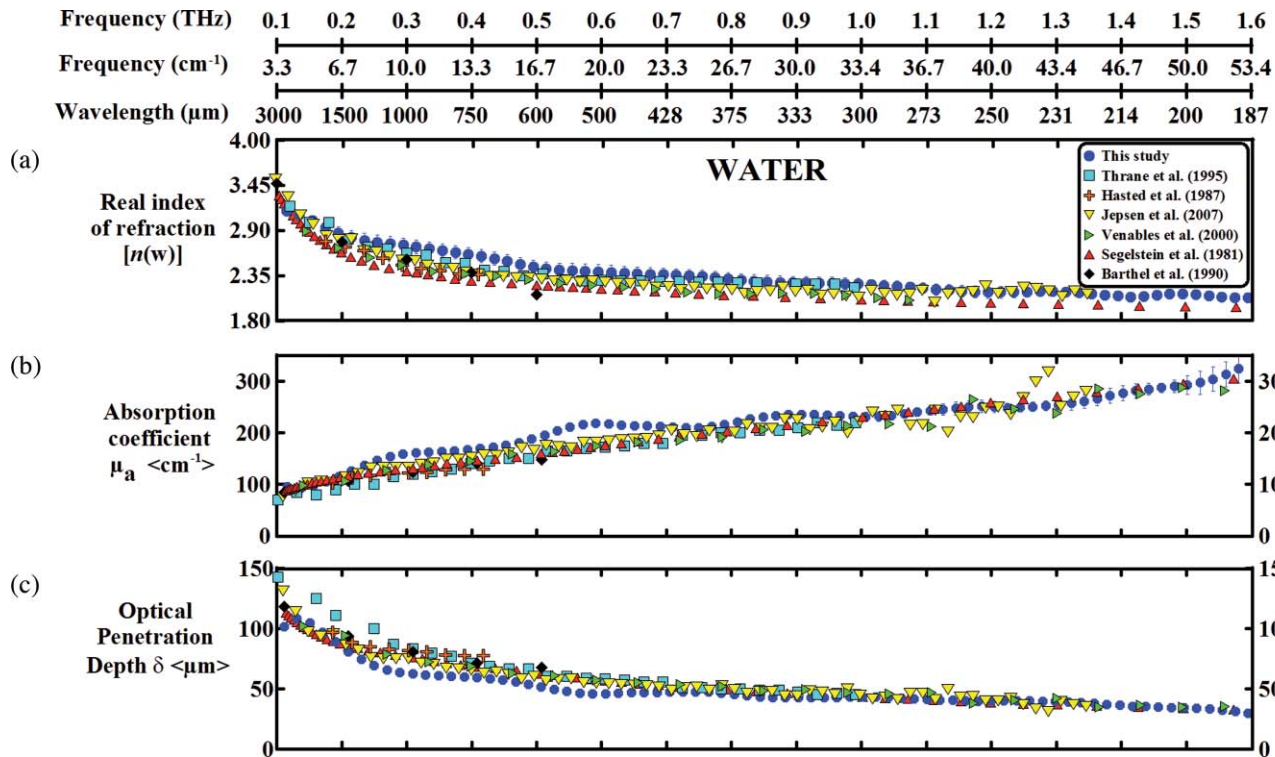
depth of roughly 1 mm. This finding suggests that ethanol does not absorb THz radiation as strongly as other polar liquids, such as water.

### 3.1.2 Water

Water was the second liquid that we used to calibrate our system. The optical properties of water, the primary constituent of all biological tissues, are also well-characterized at THz frequencies.<sup>7,28,30,32,46,48,49,66,67</sup> Figure 4(a) contains a plot of



**Fig. 3** Optical properties of ethanol at Terahertz frequencies. (a) Real index of refraction ( $n$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (b) The absorption coefficient ( $\mu_a$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (c) Optical penetration depth ( $\delta$ ) plotted versus frequency (THz). Ethanol spectra measured in this study, blue circles; fs-THz pulsed spectroscopy results from Kindt et al. (Ref. 28), purple circles; Yomogida et al. (Ref. 65), black circles; Jepsen et al. (Ref. 32), yellow circles; Barthel et al. (Ref. 64), green circles. Data are expressed as means  $\pm$  SD, with a sample population of  $n = 5$ .



**Fig. 4** Optical properties of water at Terahertz frequencies. (a) Real index of refraction ( $n$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (b) The absorption coefficient ( $\mu_a$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (c) Optical penetration depth ( $\delta$ ) plotted versus frequency (THz). Water spectra measured in this study, dark blue circles; Thrane et al. (Ref. 47), light blue squares; Hasted et al., orange crosses; Jepsen et al. (Ref. 32), yellow triangles; Segelstein et al. (Ref. 46), red triangles; Venables and Schmuttenmaer et al. (Refs. 7 and 8), green triangles. Barthel et al. (Ref. 64), black circles. Data are expressed as means  $\pm$  SD, with sample population  $n = 9$ .

the average index of refraction that we measured for water. For comparison, we also provided data plots reported by several other groups.<sup>7,28,30,32,46,49,66,67</sup> For water, we found that the magnitude of  $n$  decreases with increasing frequency and ranges between a value of 3.4 at 0.1 THz to 2.15 at 1.6 THz. Overall, the general slope and magnitude of the  $n$  values that we collected were comparable ( $\leq 8\%$ ) to those reported in previous studies.<sup>7,29,46</sup>

Figure 4(b) contains a plot of the average  $\mu_a$  values for water. Similar to ethanol, this data shows that the  $\mu_a$  of water also increases with frequency and ranges in value from 100 cm<sup>-1</sup> at 0.1 THz to 300 cm<sup>-1</sup> at 1.6 THz. For the most part, the slope and magnitude of our data is in well-agreement with previous reports.<sup>7,29,46</sup> However, since the optical absorption of water does vary with temperature,<sup>37</sup> it is possible that the slight variations between studies may be attributed to the different temperatures used in each work. Figure 4(c) shows that the  $\delta$  of water equals 125  $\mu\text{m}$  at 0.1 THz and 40  $\mu\text{m}$  at 1.6 THz.

### 3.2 Biological Tissues: Histology and Frequency Dependent Optical Properties

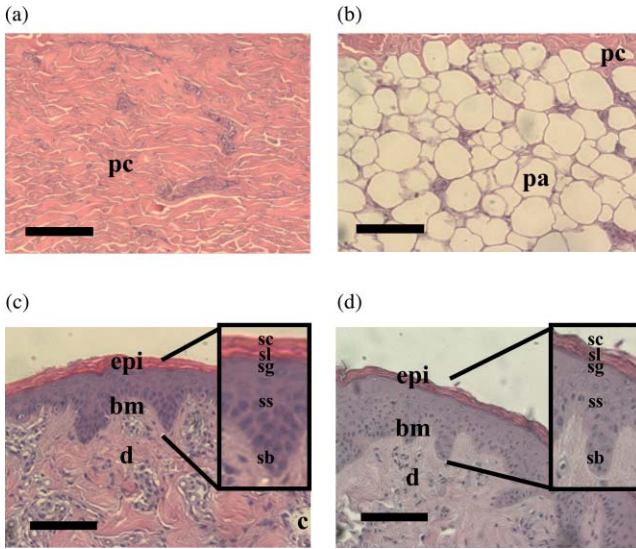
Several studies have shown that the optical properties of biological tissues vary with organism (e.g., human, rat, porcine),<sup>10,11,52,53,57</sup> tissue type (e.g., skin, muscle, adipose, liver, heart),<sup>53,58,59</sup> tissue preparation technique (e.g., freshly excised versus frozen/thawed), and duration of fixation.<sup>57</sup> Given these findings, we conducted several experiments to measure the

optical properties for skin tissue and its primary constituents, adipose tissue and muscle. In addition, to explore the effect that tissue preparation has on these values, we also collected data for both freshly excised and frozen/thawed skin tissues. The main goal for these studies was to determine whether the data collected with our compact THz-TDS system was consistent with values reported in the literature.

#### 3.2.1 Histological evaluation

Conventional histological techniques were performed to evaluate the morphology of the porcine tissues tested with our system. Immediately after THz-TDS measurements, tissues were fixed, sectioned, and stained with hematoxylin and eosin (H&E). Figures 5(a) and 5(d) contain sample representations for the following tissues: panniculus carnosus (subcutaneous striated muscle), panniculus adiposus (subcutaneous adipose), freshly harvested skin (epidermis and dermis), and frozen/thawed prepared skin (epidermis and dermis). We found that the fascicles of the striated muscle ranged in thickness between 20 and 50  $\mu\text{m}$  [Fig. 5(a)]. In addition, we found that the diameter of adipocytes ranged between 10 and 70  $\mu\text{m}$ . An image of the panniculus adiposus, a fatty layer of subcutaneous tissue superficial to the vestigial layer of muscle is provided in Fig. 5(b).

Porcine skin consists of three primary layers: epidermis, dermis, and hypodermis. In our images, the epidermis, basement membrane, and dermal tissue are all distinguishable [Figs. 5(c) and 5(d)]. In addition, in both the fresh and frozen/thawed skin



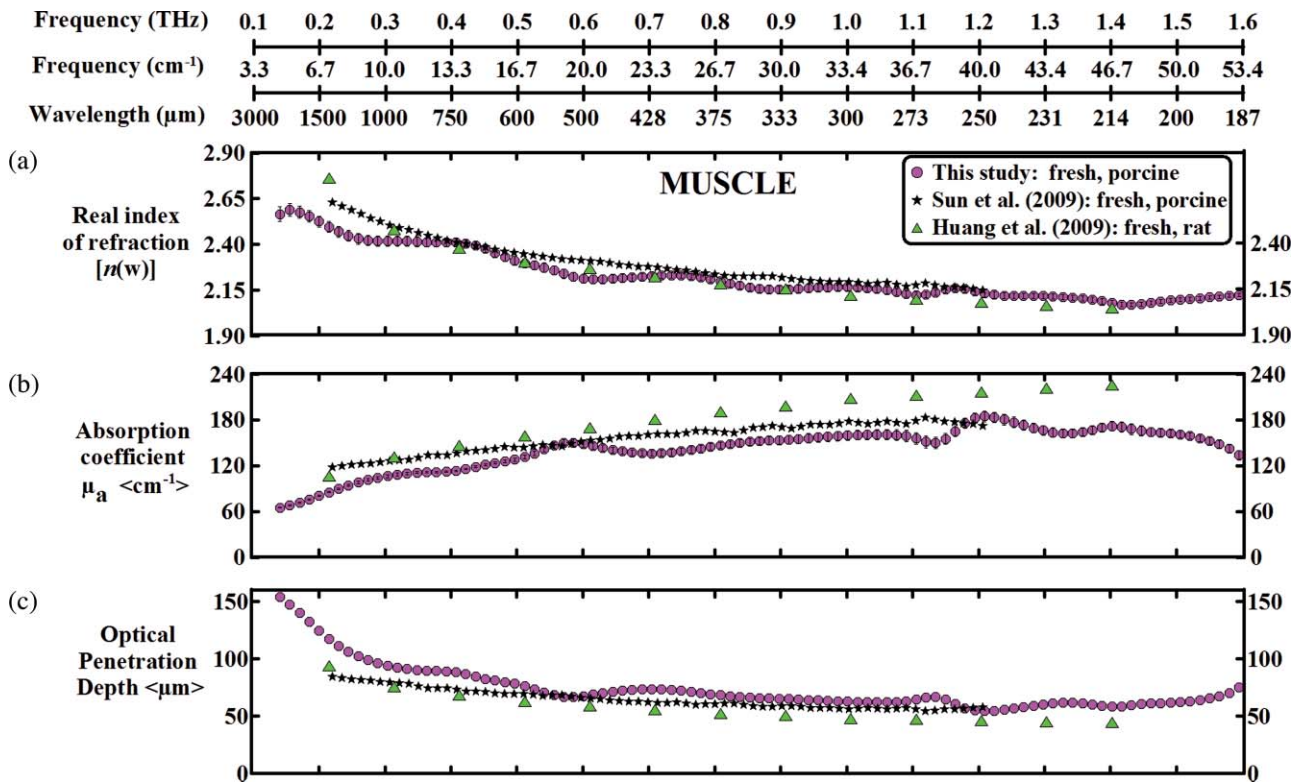
**Fig. 5** Histological evaluation of porcine cutaneous and subcutaneous tissues examined with THz-TDS system. Fixed tissue sections were stained with hematoxylin and eosin (H&E). (a) Freshly harvested muscle. (b) Freshly harvested adipose. (c) Freshly harvested skin. (d) Frozen/thawed skin. In all images the scale bar = 75  $\mu\text{m}$ . Abbreviations: muscle (panniculus carnosus = pc); adipose (panniculus adiposus = pa); epidermis = epi; basement membrane = bm; dermis = d; and capillary = c.

sections, we can differentiate all five layers of the epidermis: stratum corneum (sc), stratum lucidum (sl), stratum granulosum (sg), stratum spinosum (ss), and stratum basale (sb) [Figs. 5(c) and 5(d)]. We found that the epithelium projections ranged in thickness between 40 and 125  $\mu\text{m}$ . Given the limited penetration depth of THz radiation, these morphological features are important to note.

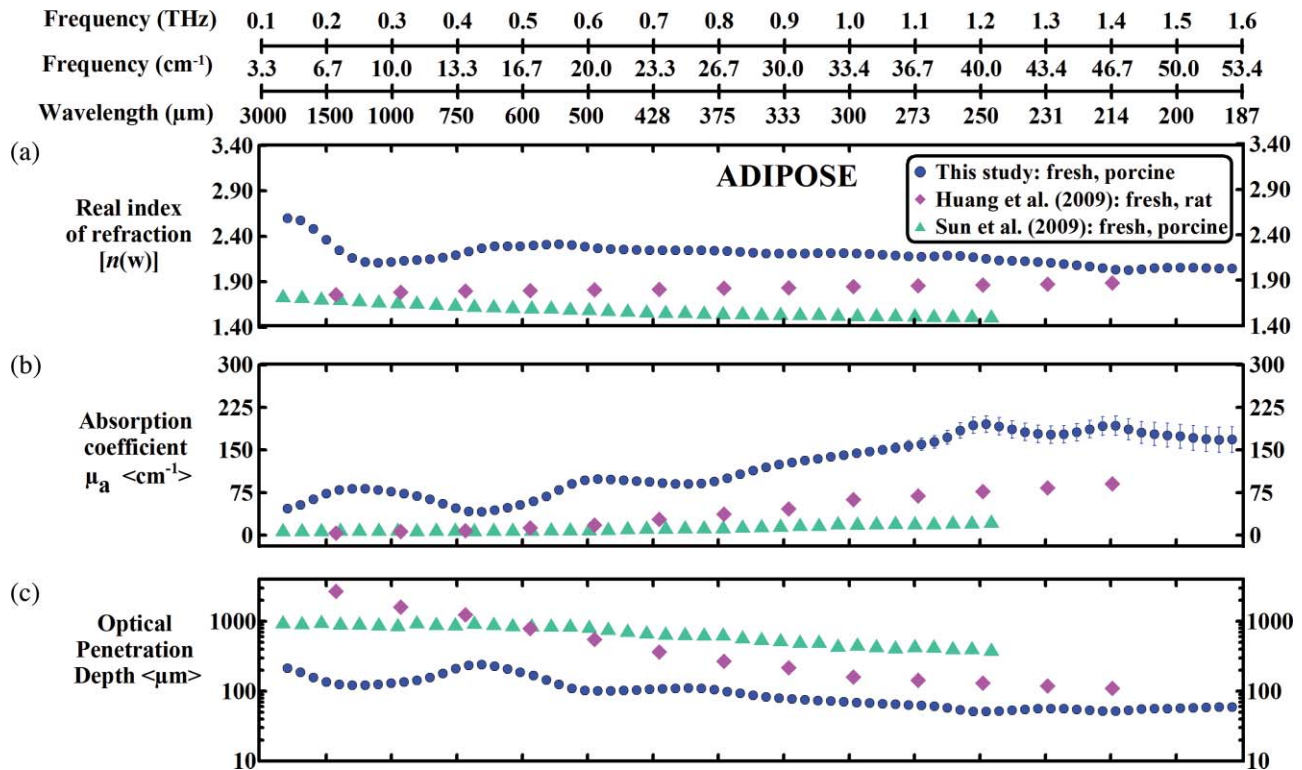
### 3.2.2 Fresh porcine muscle

Figures 6(a)–6(c) contain plots of the average optical properties and optical penetration depth for fresh porcine muscle. We found that the index of refraction decreases with frequency, whereas the absorption coefficient increases with frequency. Maximum ( $n = 2.5$ ) and minimum ( $n = 2.1$ ) values were observed at 0.13 and 1.6 THz, respectively. For comparison purposes, we also provided data plots from Sun et al.<sup>57</sup> and Huang et al.<sup>51</sup> [Figs. 6(a)–6(c)]. Overall, the  $n$  values that we measured with our system were in excellent agreement ( $\leq 10\%$ ) with those reported in both of these studies.

In contrast, we found that the  $\mu_a$  values for muscle increase with frequency, with minimum ( $\mu_a = 60 \text{ cm}^{-1}$ ) and maximum values ( $\mu_a = 200 \text{ cm}^{-1}$ ) occurring at 0.13 and 1.2 THz, respectively. These values are in excellent agreement with those reported in previous studies.<sup>51,57</sup> Figure 6(c) is a plot of the  $\delta$  of muscle at THz frequencies. The data indicates that THz radiation has a  $\delta$  in muscle of 150  $\mu\text{m}$  at 0.13 THz and 75  $\mu\text{m}$  at 1.6 THz.



**Fig. 6** Optical properties of fresh porcine muscle. (a) Real index of refraction ( $n$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (b) The absorption coefficient ( $\mu_a$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (c) Optical penetration depth ( $\delta$ ) plotted versus frequency (THz). Muscle spectra this study, pink circles; Huang et al. (Ref. 51), green triangles; Sun et al. (Ref. 57), black stars. Data are expressed as means  $\pm$  SD, with sample population = 8.



**Fig. 7** Optical properties of fresh adipose tissue. (a) Real index of refraction ( $n$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (b) The absorption coefficient ( $\mu_a$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (c) Optical penetration depth ( $\delta$ ) plotted versus frequency (THz). Adipose tissue in this study, blue circles; Sun et al. (Ref. 57), sea green triangles; Huang et al. (Ref. 51), pink diamonds. Data are expressed as means  $\pm$  SD, with sample population = 9.

It is clear from the data that the optical properties we measured were consistent with several other reports. This finding is particularly interesting because in this study we measured the properties of freshly excised porcine muscle, while other groups used were either harvested from a different organism (porcine versus rat) and/or were prepared using a different technique (fresh versus refrigerated). Specifically, Sun et al. measured the spectra of porcine muscle 24 h post-harvest, whereas Huang et al. measured spectra on freshly harvested rat muscle.<sup>52,57</sup> This data suggests that the optical properties of muscle are fairly consistent regardless of the organism or tissue preparation technique employed.

### 3.2.3 Fresh porcine adipose tissue

Figures 7(a)–7(c) contain plots of the average optical properties ( $n$  and  $\mu_a$ ) and optical penetration depth for adipose tissue. We found that the index of refraction for adipose tissue decreases with increases in frequency. Similar to muscle, maximum ( $n = 2.5$ ) and minimum ( $n = 2.0$ ) values were observed at 0.13 and 1.6 THz, respectively. The  $\mu_a$  values for adipose increases with frequency, with minimum ( $\mu_a = 50 \text{ cm}^{-1}$ ) and maximum values ( $\mu_a = 150 \text{ cm}^{-1}$ ) occurring at 0.13 and 1.6 THz, respectively. We found that the  $\delta$  in adipose ranges in value from  $200 \mu\text{m}$  at 0.13 THz to  $50 \mu\text{m}$  at 1.6 THz [Fig. 7(c)].

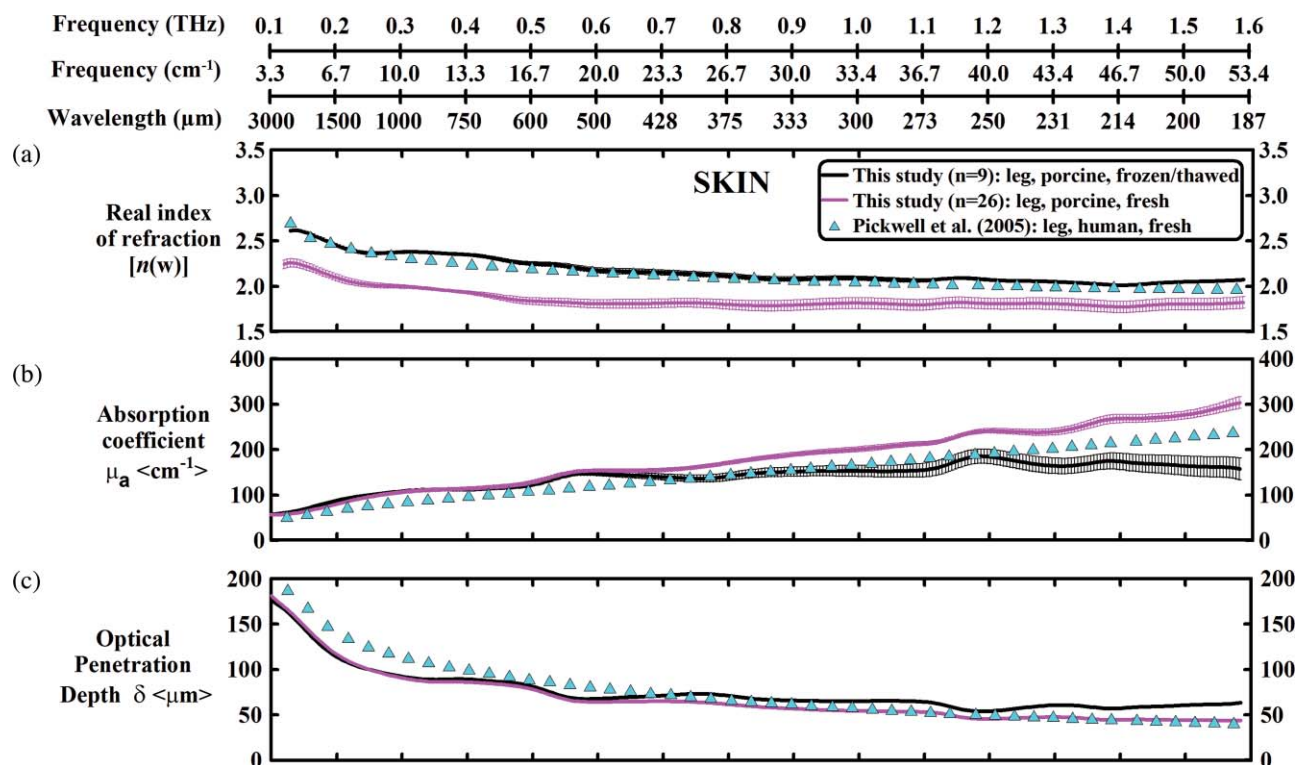
Both the slope and the index of refraction that we observed for adipose tissue were vastly different than those reported in other studies.<sup>51,57</sup> For instance, Huang et al. found that the index of refraction actually increased with frequency and ranged

in value from 1.5 at 0.2 THz to 1.9 at 1.4 THz. In addition, Sun et al. found that the index of refraction decreased with frequency, and ranged in value from 1.7 at 0.15 THz to  $\sim 1.4$  at 1.2 THz. Overall, these values are lower ( $\sim 30\%$ ) than the index of refraction values that we measured with our system. Additionally, we found that the absorption coefficient values that we measured for adipose tissue were roughly two to ten times greater in magnitude than those previously reported.<sup>51,57</sup> Although the exact reason for these differences is unclear, it is possible that during our measurements a thin layer of water was present on the adipose tissue. This may also explain why the properties that we measured for adipose are similar to those that we measured for water. In contrast, Huang and Sun et al. used blotting paper or ethanol to further dehydrate their adipose tissues prior to measurement. Such approaches are advantageous for collecting spectra from adipose tissue, but it is also important to keep in mind that in an *in vivo* scenario adipose tissue will be hydrated. Thus, although our data is inconsistent with previous reports, it may be more indicative of the values that may be measured on hydrated, intact biological tissues. Nevertheless, the differences in optical properties that we observed for adipose tissue are most likely due to the fact that we used different sample preparation and dehydration techniques.

### 3.2.4 Measurement and comparison of the optical properties of skin

After measuring the spectra for muscle and adipose tissue, we then collected and compared the properties for fresh and





**Fig. 8** Optical properties of excised porcine skin prepared using fresh and frozen/thawed techniques. (a) Real index of refraction ( $n$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (b) The absorption coefficient ( $\mu_a$ ) plotted versus frequency (THz) and wavelength ( $\mu\text{m}$ ). (c) Optical penetration depth ( $\delta$ ) plotted versus frequency (THz). Spectra for fresh excised porcine skin from left leg, pink line; frozen/thawed skin from leg, black line; human skin data from Pickwell et al. (Ref. 10), light blue triangles. Data are expressed as means  $\pm$  SD, with  $n = 9$  or 26.

frozen/thawed skin tissues. Plots for the index of refraction, absorption coefficient, and optical penetration depth are provided in Figs. 8(a)–8(c). Similar to previous tissues, we found that the index of refraction decreases with increases in frequency. For fresh skin, we found that the  $n$  ranged between 2.3 and 1.8, whereas for the frozen/thawed samples the  $n$  values were slightly higher and ranged in value between 2.6 and 2.0. The values we measured for the frozen/thawed samples were in excellent agreement with those reported by Pickwell et al.<sup>53</sup>

Similar to the trends observed in other tissues, we also found that the magnitude of  $\mu_a$  increased with frequency. At lower THz frequencies (0.1 to 0.8 THz), the fresh and frozen/thawed skin samples exhibited comparable values for  $\mu_a$  (50 to 180 cm<sup>-1</sup>), which were both in excellent agreement with those values measured by Pickwell et al. In contrast, for frequencies greater than 0.8 THz, we found that fresh tissue exhibited  $\mu_a$  values that were 10 to 100% greater than those measured for frozen/thawed tissues. Interestingly, Pickwell et al. reported  $\mu_a$  values that fell between these values. Figure 8(c) is a plot of the  $\delta$  of THz radiation in fresh excised porcine skin. The data shows that THz radiation will have maximal penetration ( $\delta \approx 200 \mu\text{m}$ ) at lower frequencies and minimal penetration at higher THz frequencies ( $\delta \approx 40$  to  $50 \mu\text{m}$ ). These penetration depths are important to consider given the fact that the thickness of the epidermis ranges across the human body between 50 and 1000  $\mu\text{m}$ . In addition, we found that the epithelial thickness for the same section of porcine skin ranged in thickness between 40 and 125  $\mu\text{m}$  [Figs. 5(c) and 5(d)]. Since the optical penetration depth of THz radiation in

skin is comparable to these distances, it is entirely possible that THz-TDS measurements may fluctuate with the thickness of the epithelium projections that are present in the sampled area.

## 4 Conclusion

In summary, we have developed and tested the performance of a compact THz-TDS device that is considerably smaller and lighter than conventional systems. Using this device, we measured the optical properties from 0.1 to 1.6 THz for ethanol, water, muscle, adipose tissue, and skin. In addition, we compared the spectra that we collected to values reported in the literature. Consistent with previously published reports, we found that the index of refraction decreases with frequency for all liquids and tissues tested. Muscle, adipose, and frozen/thawed skin exhibited comparable  $n$  values ranging from 2.5 at 0.1 THz to 2.0 at 2.0 THz; however, fresh skin exhibited values that were roughly 40% lower in magnitude. Also consistent with previous reports, we found that the  $\mu_a$  measured for each sample increases with frequency. Muscle and adipose tissue exhibited  $\mu_a$  values ranging between 50 and 200 cm<sup>-1</sup>, whereas both fresh and frozen/thawed skin samples exhibited slightly higher  $\mu_a$  values ranging between 50 and 300 cm<sup>-1</sup>. Additionally, we determined that the  $\delta$  in skin was roughly 150  $\mu\text{m}$  at 0.1 THz and 40 to 50  $\mu\text{m}$  at 1.6 THz. Overall, our compact THz-TDS system provided measures for the optical properties of biological tissues that were consistent with previous studies. These results demonstrate that our compact THz spectrometer may be a

useful and practical tool for the rapid battlefield assessment of skin health, wounds, and burns.

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